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# The Crystal Structure of Sodium Pyruvate* 

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Sodium pyruvate, $\mathrm{CH}_{3} \mathrm{COCOONa}$, crystallized from aqueous solution, is monoclinic, space group $P 2_{1} / a$, with four molecules per unit cell of dimensions

$$
a=22.25, b=5.31, c=3.71 \AA ; \beta=98.2^{\circ} .
$$


#### Abstract

The detailed structure has been determined by two-dimensional Fourier syntheses, and refined by the method of least-squares.

The plane containing the methyl carbon, the keto group and the carbon of the carboxylic group makes an angle of $18 \cdot 1^{\circ}$ with the plane of the $\mathrm{C}-\mathrm{COO}$ group. This suggests that there is no appreciable degree of conjugation across the central $\mathrm{C}-\mathrm{C}$ bond, a fact that is confirmed by the observed central C-C bond length, $1.579 \pm 0.045 \AA$, which is close to the standard single $\mathrm{C}-\mathrm{C}$ bond length.

Strong $\mathrm{Na}-\mathrm{O}$ bonds tie the molecules into infinite layers parallel to the (100) planes. Any one layer of molecules is linked with the neighbouring layers by strong $\mathrm{Na}-\mathrm{O}$ bonds on one side and van der Waals bonds on the other.


## Introduction

The $\alpha$-keto acid analogs of amino acids are of considerable biochemical interest as intermediates in the biosynthesis and degradation of amino acids. They have unusual properties and are formed in the course of enzymatic oxidative deamination and transamination reactions. Although considerable work has been reported on the structure of amino acids, little attention has been given so far to the $\alpha$-keto analogs. With a view to providing structural information which might throw light on their chemical characteristics, we have undertaken a programme to determine the structures of some of these compounds. In what follows, we report the structure analysis of the sodium salt of the simplest keto acid, i.e. pyruvic acid.

## Experimental details

The crystals of sodium pyruvate were obtained from aqueous solution. They are monoclinic and grow as very thin plates, parallel to the (100) face. There is strong cleavage parallel to the plates. The crystal and physical data obtained are as follows:

$$
a=22 \cdot 25, b=5 \cdot 31, c=3 \cdot 71 \AA ; \beta=98 \cdot 2^{\circ} .
$$

The axial lengths were determined from the rotation

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photographs along the three axes, and the $\beta$ angle from a Laue photograph, taken with X-rays, travelling parallel to the $b$ axis.

The observed density at $25^{\circ} \mathrm{C}$., measured by the flotation method is 1.718 g.cm..$^{-3}$, and the calculated density for four molecules of $\mathrm{CH}_{3} \mathrm{COCOONa}$ per unit cell is $\mathrm{I} \cdot 684$ g.cm..$^{-3}$.

The systematic absences are: $0 k 0$ for $k$ odd and $h 0 l$ for $h$ odd, so that the space group is uniquely determined as $P 2_{1} / a$. The linear absorption coefficient $\mu$ for $\mathrm{Cu} K \alpha$ radiation is $23.7 \mathrm{~cm} .^{-1}$. Reflexions of the type $h k 0, h 0 l$ and $0 k l$ were obtained from $10^{\circ}$ or $15^{\circ}$ oscillation photographs, using $\mathrm{Cu} K \alpha$ radiation and the multiple-film technique. Of the 153 possible reflexions in the $h k 0$ zone, 122 were observed while in the $h 0 l$ and $0 k l$ zones, 93 and 19 reflexions were observed out of possible 105 and 27 , respectively. The specimens used were flakes of roughly ( $0 \cdot 7 \times 1 \cdot 1$ ) mm. ${ }^{2}$ area. Intensities were measured visually, using intensity scales made with the same crystals. Corrections for Lorentz and polarization factors were applied and that for absorption was neglected. On account of the flakiness of the crystals, the shapes and sizes of the spots were not uniform, and this seems to be a serious source of error in the estimation of integrated intensities.

Wilson's method (1942) was used to obtain preliminary values for the scale factors and the temperature factors. These values were improved during the later
stages of refinement. The final overall temperature factors for the three zones are: $1 \cdot 1 \AA^{2}$ for $F(h k 0)$, $2 \cdot 1 \AA^{2}$ for $F(h 0 l)$ and $3 \cdot 2 \AA^{2}$ for $F(0 k l)$. Final structure factors were calculated using McWeeny's (195l) values of atomic scattering factors for carbon and oxygen and James \& Brindley's values for $\mathrm{Na}^{+}$.

## Determination and refinement of the structure

The $x$ and $y$ coordinates of the sodium atom were obtained from the (001) Patterson projection. Signs of $F(h k 0)$ were calculated on the basis of sodium positions alone, and an electron-density projection obtained. Assuming usual bond lengths and angles and a plane molecule nearly parallel to the plane of projection, a trial structure was then postulated and set up. This structure was refined by the usual iterative process till the structure factors stopped changing signs, except for a few small ones. At this stage, $R$ was about 0.25 . The coordinates were then refined twice by the method of least-squares. The maximum shift in coordinates after the final refinement was $0.03 \AA$. Four reflexions, whose $F_{o}$ values were considered uncertain were not used in the final refinement. The final value of $R$ was $0 \cdot 220$ on including all except very-highangle unobserved reflexions at half the minimum observable value. On omitting the 200 reflexion, which appears to be subject to extinction, $R$ decreased to $0 \cdot 215$. The ( 001 ) electron-density projection, shown in Fig. 1, was obtained using final calculated signs, and $F_{o}$ values for all except 200,110 and 510 reflexions. For these three reflexions, final $F_{c}$ values were used.


Fig. 1. Electron density projected on (001). Contours drawn at intervals of 1 e. $\AA^{-2}$, starting from 2 e. $\AA^{-2}$.

The (010) Patterson projection gave the approximate $z$ coordinates of the sodium atom and the carboxylic oxygen atoms. The ( 010 ) electron-density projection was plotted on the basis of these atomic positions alone, and was refined by the usual method till the structure factors stopped changing signs, except
for a few small ones. At this stage, $R$ was about 0.3 . The $z$ coordinates were refined twice by the method of least-squares. Few coordinates, which were showing large shifts, were refined further till the maximum shift in coordinates was less than $0.02 \AA$. The final value of $R$ was $0 \cdot 223$, which dropped to $0 \cdot 192$ on omitting $200,201,20 \overline{1}, 10 \overline{2}$ and 202 reflexions. These five reflexions appear to be subject to extinction, and were not used in later refinements. The (010) electrondensity projection, shown in Fig. 2, was obtained using final calculated signs, and $F_{o}$ values for all except the five reflexions mentioned above, for which final $F_{c}$ values were used.


Fig. 2. Electron density projected on (010). Contours drawn at intervals of $1 \mathrm{e} . \AA^{-2}$, starting from 2 e. $\AA^{-2}$.

The $R$ factor for $0 k l$ reflexions was $0 \cdot 210$. The final atomic parameters are listed in Table 1, and the observed and calculated structure factors in Table 2.

Table 1. Final atomic parameters

|  | $x / a$ | $y / b$ | $z / c$ |
| :--- | :---: | :---: | :---: |
| Na | $0.1989 \pm 0.0008$ | $0.1802 \pm 0.0032$ | $0.1638 \pm 0.0054$ |
| $\mathrm{O}_{1}$ | $0.2035 \pm 0.0010$ | $0.4887 \pm 0.0043$ | $0.6491 \pm 0.0073$ |
| $\mathrm{O}_{2}$ | $0.1004 \pm 0.0010$ | $0.4100 \pm 0.0043$ | $0.1617 \pm 0.0073$ |
| $\mathrm{O}_{3}$ | $0.1707 \pm 0.0010$ | $0.8843 \pm 0.0043$ | $0.7144 \pm 0.0073$ |
| $\mathrm{C}_{1}$ | $0.1657 \pm 0.0013$ | $0.6608 \pm 0.0057$ | $0.5854 \pm 0.0097$ |
| $\mathrm{C}_{2}$ | $0.1024 \pm 0.0013$ | $0.5803 \pm 0.0057$ | $0.3685 \pm 0.0097$ |
| $\mathrm{C}_{3}$ | $0.0523 \pm 0.0013$ | $0.7722 \pm 0.0057$ | $0.3882 \pm 0.0097$ |

## Estimation of errors

The standard deviations of atomic coordinates and the interatomic distances were estimated (Lipson \& Cochran, 1953) assuming $\sigma(F)=0 \cdot 2|F|$, and $p=5 \AA-2$. A root mean square of $\sigma(x), \sigma(y)$ and $\sigma(z)$ was taken as the standard deviation of each atomic position.

The standard deviations in bond angles were estimated by the method of Ahmed \& Cruickshank (1953).

## Description of the structure and discussion

## (a) Molecular structure

The bond lengths and bond angles found in the pyruvate group are shown in Fig. 3 and listed in Table 3. The mean plane through the atoms $\mathrm{C}_{2}, \mathrm{C}_{1}$, $\mathrm{O}_{1}$ and $\mathrm{O}_{3}$ can be represented by the equation

Table 2. Observed and calculated structure factors


$$
1 \cdot 6813 x^{\prime}+1 \cdot 2912 y-4 \cdot 3267 z^{\prime}-3 \cdot 4032=0
$$

where $x^{\prime}, y$ and $z^{\prime}$ are expressed in $\AA$ units and referred to the orthogonal axes $a^{\prime}, b$ and $c$. The perpendicular distances of the individual atoms from the mean plane are $0.029 \AA$ for $\mathrm{O}_{1}, \mathrm{C}_{2}$ and $\mathrm{O}_{3}$ and $-0.029 \AA$ for $\mathrm{C}_{1}$.


Fig. 3. Bond lengths and bond angles in the pyruvate group. $\mathrm{O}_{1}$ and $\mathrm{O}_{3}$ are carboxylic oxygens, $\mathrm{O}_{2}$ is keto group oxygen and $\mathrm{C}_{3}$ is methyl carbon.

Referred to the same axes $a^{\prime}, b$ and $c$, the mean plane through the atoms $\mathrm{C}_{1}, \mathrm{C}_{2}, \mathrm{O}_{2}$ and $\mathrm{C}_{3}$ can be represented by the equation

$$
1 \cdot 3021 x^{\prime}+2 \cdot 8726 y-4 \cdot 1758 z^{\prime}-7 \cdot 6983=0
$$

The perpendicular distances of the individual atoms from the mean plane are $0.050 \AA$ for $\mathrm{O}_{2}, \mathrm{C}_{1}$ and $\mathrm{C}_{3}$ and $-0.050 \AA$ for $\mathrm{C}_{2}$.

The two planes make an angle of $18 \cdot 1^{\circ}$ with each other.

Table 3. Intramolecular bond lengths and angles

|  | Bond length |  | Bond angles |
| :--- | :--- | :--- | :--- |
| $\mathrm{C}_{1}-\mathrm{O}_{1}$ | $1 \cdot 242 \pm 0 \cdot 040 \AA$ | $\mathrm{O}_{1}-\mathrm{C}_{1}-\mathrm{O}_{3}$ | $126 \cdot 3 \pm 2 \cdot 5^{\circ}$ |
| $\mathrm{C}_{1}-\mathrm{O}_{3}$ | $1 \cdot 278 \pm 0 \cdot 040$ | $\mathrm{O}_{1}-\mathrm{C}_{1}-\mathrm{C}_{2}$ | $115 \cdot 4 \pm 2 \cdot 4$ |
| $\mathrm{C}_{1}-\mathrm{C}_{2}$ | $1 \cdot 579 \pm 0 \cdot 045$ | $\mathrm{O}_{3}-\mathrm{C}_{1}-\mathrm{C}_{2}$ | $117 \cdot 8 \pm 2 \cdot 4$ |
| $\mathrm{C}_{2}-\mathrm{O}_{2}$ | $1 \cdot 183 \pm 0 \cdot 040$ | $\mathrm{C}_{1}-\mathrm{C}_{2} \mathrm{O}_{2}$ | $119 \cdot 1 \pm 2 \cdot 5$ |
| $\mathrm{C}_{2}-\mathrm{C}_{3}$ | $1 \cdot 518 \pm 0 \cdot 045$ | $\mathrm{C}_{1}-\mathrm{C}_{2}-\mathrm{C}_{3}$ | $113 \cdot 7 \pm 2 \cdot 3$ |
|  |  | $\mathrm{O}_{2}-\mathrm{C}_{2}-\mathrm{C}_{3}$ | $125 \cdot 8 \pm 2 \cdot 6$ |

The central $\mathrm{C}_{1}-\mathrm{C}_{2}$ bond length ( $1.579 \pm 0.045 \AA$ ) is not significantly different from the normal single C-C bond length. This suggests that there is no appreciable


Fig. 4. Structure projected on (001).


Fig. 5. Structure projected on (010).
degree of conjugation across this bond. There would thus be a possibility of free rotation of the two planes across the central $\mathrm{C}-\mathrm{C}$ bond. This explains the fact that the two planes $\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{O}_{2} \mathrm{C}_{3}$ and $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{O}_{3} \mathrm{O}_{1}$ are not coplanar but make an angle of $18 \cdot 1^{\circ}$ with each other.

The two bonds $\mathrm{C}_{1}-\mathrm{O}_{1}$ and $\mathrm{C}_{1}-\mathrm{O}_{3}$ are of different lengths, although the difference may be within experimental error. In oxalic acid and eight aminoacids listed by Ahmed \& Cruickshank (1953), the two $\mathrm{C}-\mathrm{O}$ bonds are always of unequal lengths.

The angle $\mathrm{O}_{1}-\mathrm{C}_{1}-\mathrm{O}_{3}\left(126.3 \pm 2.5^{\circ}\right)$ is significantly larger than the angles $\mathrm{C}_{2}-\mathrm{C}_{1}-\mathrm{O}_{1}\left(115 \cdot 4 \pm 2 \cdot 4^{\circ}\right)$ and $\mathrm{C}_{2}-\mathrm{C}_{1}-\mathrm{O}_{3}\left(117 \cdot 8 \pm 2 \cdot 4^{\circ}\right)$. This is in agreement with the results found for oxalic acid and the amino-acids listed by Ahmed \& Cruickshank (1953).

The differences in the angles $\mathrm{O}_{2}-\mathrm{C}_{2}-\mathrm{C}_{3}\left(125 \cdot 8 \pm 2 \cdot 6^{\circ}\right)$, $\mathrm{O}_{2}-\mathrm{C}_{2}-\mathrm{C}_{1}\left(119 \cdot 1 \pm 2 \cdot 5^{\circ}\right)$ and $\mathrm{C}_{3}-\mathrm{C}_{2}-\mathrm{C}_{1}\left(113 \cdot 7 \pm 2 \cdot 3^{\circ}\right)$ are also significant.

On account of the rather large uncertainties in the observed bond lengths and angles, a discussion of the molecular structure with reference to its chemical properties or to the molecular structures of related compounds is not warranted at this stage.

## (b) Crystal structure

The projections of the structure along the $c$ and $b$ axes are shown in Fig. 4 and 5 respectively.

There are six short bonds and one slightly longer bond between $\mathrm{Na}(x, y, z)$ and the neighbouring oxygens; these distances are as follows:

$$
\begin{array}{ll}
\mathrm{Na}-\mathrm{O}_{1}(x, y, z) & 2 \cdot 43 \pm 0.03 \AA \\
\mathrm{Na}-\mathrm{O}_{1}(x, y, z-1) & 2 \cdot 63 \pm 0.03 \\
\mathrm{Na}-\mathrm{O}_{3}(x, y-1, z) & 2 \cdot 72 \pm 0.03 \\
\mathrm{Na}-\mathrm{O}_{3}(x, y-1, z-1) & 2 \cdot 31 \pm 0.03 \\
\mathrm{Na}-\mathrm{O}_{1}\left(\frac{1}{2}-x, y-\frac{1}{2}, 1-z\right) & 2 \cdot 41 \pm 0.03 \\
\mathrm{Na}-\mathrm{O}_{2}(x, y, z) & 2 \cdot 51 \pm 0.03 \\
\mathrm{Na}-\mathrm{O}_{3}\left(\frac{1}{2}-x, y-\frac{1}{2}, 1-z\right) & 3 \cdot 07 \pm 0.03
\end{array}
$$

There are van der Waals bonds between the methyl carbons $\mathrm{C}_{3}$ and keto group oxygens $\mathrm{O}_{2}$, and between carbons and carbons of the neighbouring molecules. The more significant distances are as follows:

| $\mathrm{C}_{3}(\bar{x}, 1-y, \bar{z})-\mathrm{C}_{3}(x, y, z)$ | $4 \cdot 49 \pm 0.06 \AA$ |
| :--- | :--- |
| $\mathrm{C}_{3}(\bar{x}, 1-y, \bar{z})-\mathrm{C}_{3}(x, y, z-1)$ | $3 \cdot 87 \pm 0 \cdot 06$ |
| $\mathrm{C}_{3}(\bar{x}, 1-y, \bar{z})-\mathrm{C}_{3}(x, y-1, z)$ | $4 \cdot 20 \pm 0 \cdot 06$ |
| $\mathrm{C}_{3}(\bar{x}, 1-y, \bar{z})-\mathrm{C}_{3}(x, y-1, z-1)$ | $3 \cdot 53 \pm 0 \cdot 06$ |
| $\mathrm{C}_{3}(\bar{x}, 1-y, \bar{z})-\mathrm{O}_{2}(x, y, z-1)$ | $4 \cdot 11 \pm 0 \cdot 04$ |
| $\mathrm{C}_{3}(\bar{x}, \mathrm{I}-y, \bar{z})-\mathrm{O}_{2}(x, y, z)$ | $3 \cdot 83 \pm 0.04$ |

The strong $\mathrm{Na}-\mathrm{O}$ bonds tie the molecules into infinite layers parallel to the (100) plane. Anyone layer of molecules is linked with the neighbouring layers tightly on one side by bonds such as $\mathrm{Na}(x, y, z)-\mathrm{O}_{1}\left(\frac{1}{2}-x, y-\frac{1}{2}, 1-z\right)$, and loosely on the other side by van der Waals bonds, mentioned above. This arrangement explains the presence of strong cleavage parallel to the (100) face.

The refinement of the temperature factors has not been attempted. The differences in the overall temperature factors for the three projections, however,
suggest that the mean amplitude of the thermal vibration of atoms has the largest component along the $c$ axis and the smallest component along the $a$ axis, since the overall temperature factor is least for the ( 001 ) projection ( $1 \cdot 1 \AA^{2}$ ) and maximum for the (100) projection ( $3 \cdot 2 \AA^{2}$ ). This is also indicated by the (010) electron-density projection (Fig. 2) where the atomic contours, especially those of the methyl carbon atom, appear distinctly elongated nearly along the $c$ axis. This can be understood from the orientation of the molecule ( $x, y, z$ ) with respect to the neighbouring molecules in Fig. 4. The atom $\mathrm{O}_{1}$ is strongly bonded to the two sodium atoms $(x, y, z)$ and $(x, y, z+1)$, and atom $\mathrm{O}_{3}$ is strongly bonded to the two sodium atoms $(x, y+1, z)$ and $(x, y+1, z+1)$. The methyl carbon $\mathrm{C}_{3}$ and keto-group oxygen $\mathrm{O}_{2}$ are bonded to the neighbouring molecules by weak forces only. The oscillation of the molecule about the line $\mathrm{O}_{1}-\mathrm{O}_{3}$ should therefore be most likely, and so the components of the thermal
vibration amplitudes of the atoms $\mathrm{C}_{2}, \mathrm{C}_{3}$ and $\mathrm{O}_{2}$ will be maximum along the $c$ axis and will be least along the $a$ axis. The bond between $\mathrm{Na}(x, y, z)$ and $\mathrm{O}_{2}(x, y, z)$ is quite strong (distance $-2.51 \AA$ ), but since the movement of the atom $\mathrm{O}_{2}$ would take place nearly perpendicular to the $\mathrm{Na}-\mathrm{O}_{2}$ bond, the $\mathrm{Na}-\mathrm{O}_{2}$ distance is not expected to alter appreciably during the oscillation. Besides, there is a possibility of free rotation about the central $\mathrm{C}-\mathrm{C}$ bond, so that the plane $\mathrm{C}_{3}-\mathrm{C}_{2}-\mathrm{O}_{2}$ may keep turning suitably during the oscillation so that the $\mathrm{Na}-\mathrm{O}_{2}$ distance is not affected.

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## Short Communications

Contributions intended for publication under this heading should be expressly so marked; they should not exceed about 1000 words; they should be forwarded in the usual way to the appropriate Co-editor; they will be published as speedily as possible. Publication will be quicker if the contributions are without illustrations.

Acta Cryst. (1961). 14, 1286
A note on the crystal structure of zirconium pyrophosphate. By H. McD. McGeachin, Albright and Wilson (Mfg) Ltd., Oldbury, Birmingham, England
(Received 26 June 1961)

Levi \& Peyronel's (1935) structure of zirconium pyrophosphate $\mathrm{ZrP}_{2} \mathrm{O}_{7}$ requires that the central $\mathrm{P}-\mathrm{O}-\mathrm{P}$ group of the pyrophosphate ion be linear, lying in fact on a three-fold axis. In both paper and abstract there occurs an unfortunate numerical error by which the P-P distance is incorrectly given as $3.03 \AA$ which is less than twice $1.56 \AA$, the P-O distance. This has led some writers (Hanwick \& Hoffmann, 1951; Van Wazer, 1958) to the conclusion that the $\mathrm{P}-\mathrm{O}-\mathrm{P}$ angle is $152^{\circ}$ and not $180^{\circ}$. Furthermore it has been felt that since the X-ray scattering is dominated by the zirconium atoms, the determination of the structure of the pyrophosphate ion may not be unequivocal.

The low temperature form of $\mathrm{ZrP}_{2} \mathrm{O}_{7}$ crystallizes in a primitive cubic cell of side $a_{0}=8 \cdot 258 \AA$ (Harrison, McKinstry \& Hummel, 1954). Levi \& Peyronel give the density as about 3.3 g.cm..$^{-3}$, which corresponds to four formula units per cell. In any such cell at least some of the eight phosphorus atoms must lie on three-fold axes. If $\mathrm{P}_{2} \mathrm{O}_{7}^{4-}$ is a discrete ion and one phosphorus atom lies on a three-fold axis, so must the other, for otherwise the ion would contain four phosphorus atoms; the only sensible way of distributing the seven oxygen atoms is to place one on the axis and six in two symmetrical
groups of three. To preserve four-fold grouping of oxygen atoms round the phosphorus atoms the only possible configuration is one in which the single oxygen atom is placed on the three-fold axis between the phosphorus atoms.
The existence and linearity of the central $\mathrm{P}-\mathrm{O}-\mathrm{P}$ group in $\mathrm{ZrP}_{2} \mathrm{O}_{7}$ can thus be confirmed (without detailed knowledge of the structure or even the space group) from considerations of the Bravais lattice, cell size and density, and the assumption that the $\mathrm{P}_{2} \mathrm{O}_{7}$ ion is discrete. It should be noted, however, that in $\mathrm{Na}_{4} \mathrm{P}_{2} \mathrm{O}_{7} \cdot 10 \mathrm{H}_{2} \mathrm{O}$ the P-O-P angle is $134^{\circ}$ (MacArthur \& Beevers, 1957).

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